

# Triplet state photophysics in an arylenethynylene $\pi$ -conjugated polymer

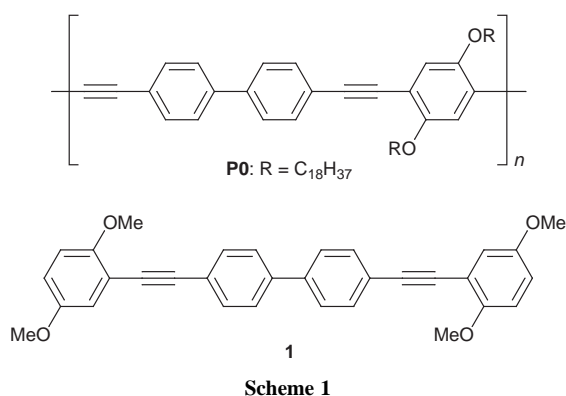
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The triplet state photophysics of a polyaryleneethynylene  $\pi$ -conjugated polymer and a related model compound are investigated with various techniques, including laser flash photolysis and time-resolved thermal lensing (TRTL).

The photophysical and photochemical properties of  $\pi$ -conjugated polymers have been of significant recent interest.<sup>1</sup> Although most work on the photophysics of  $\pi$ -conjugated polymers has focused on the  $^1\pi,\pi^*$  manifold, considerable evidence suggests that  $^3\pi,\pi^*$  excitons are also involved in their rich photochemistry.<sup>2–5</sup> Unfortunately, owing to the absence of phosphorescence in most  $\pi$ -conjugated polymers<sup>2</sup> little is known about the absolute energies and yields of the triplets produced by direct excitation.

We recently initiated a photophysical study of arylenethynylene-based  $\pi$ -conjugated polymers that contain a photoactive transition metal chromophore in the  $\pi$ -conjugated backbone.<sup>6,7</sup> A primary objective of this work is to understand the interaction between backbone-based  $^1\pi,\pi^*$  and  $^3\pi,\pi^*$  excited states and metal complex-based  $d \rightarrow \pi^*$  charge transfer excitations. However, in order to pursue this line of investigation in the metal-organic polymers, it is necessary to have a firm understanding of the properties of the  $^1\pi,\pi^*$  and  $^3\pi,\pi^*$  excited states of the all-organic polymers. Here, we report the results of a detailed photophysical study which compares the photophysical properties of the singlet and triplet manifolds of arylenethynylene polymer **P0** and model compound **1** in dilute THF solution (Scheme 1).<sup>‡</sup>



Pulsed laser excitation (355 nm, 10 ns FWHM, 4 mJ pulse<sup>-1</sup>) of **1** and **P0** affords a long lived transient [ $\tau(\mathbf{1}) = 142 \mu\text{s}$ ,  $\tau(\mathbf{P0}) = 169 \mu\text{s}$ ], which is assigned to a  $^3\pi,\pi^*$  state. The triplet state difference absorption spectra of **1** and **P0** (Fig. 1) are characterized by bleaching of the ground state  $\pi,\pi^*$  absorption band [ $\lambda_{\text{max}}(\mathbf{1}) = 360 \text{ nm}$ ,  $\lambda_{\text{max}}(\mathbf{P0}) = 410 \text{ nm}$ ] and a broad (triplet-triplet) absorption band that extends to longer wavelengths. Interestingly, the difference spectra of **1** and **P0** are remarkably similar, except that the bleach and absorption bands are red-shifted in the polymer. The qualitative similarity in the triplet-triplet absorption spectra implies that the electronic structure of the triplet state is similar in the model and polymer.

In order to determine the triplet energies of **1** and **P0** ( $E_T$ ), Stern-Volmer quenching studies were carried out with a series

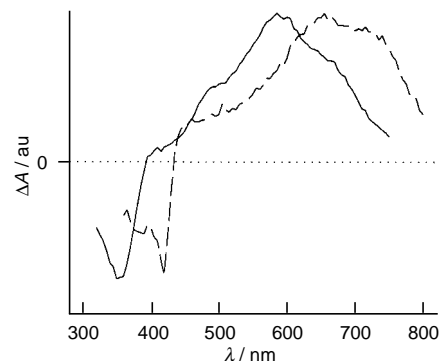


Fig. 1 Triplet state difference absorption spectra of (solid line) **1** and (dashed line) **P0** in degassed THF solution

of triplet acceptors of varying energy.<sup>8,9</sup> Triplet quenching rate constants ( $k_q$ ) were obtained from Stern-Volmer analysis of the observed triplet decay lifetimes in vacuum degassed THF solutions as a function of quencher concentration. The triplet energies were then estimated by fitting the experimental Sandros plots (Fig. 2)<sup>10</sup> with a Marcus equation that is appropriate for bimolecular reactions which occur at or near diffusion control,<sup>11,12</sup>

$$k_q = \frac{k_d}{1 + \exp(\Delta G/RT) + (k_{-d}/k_{\text{en}}^0) \exp(\Delta G^\ddagger)} \quad (1a)$$

$$\Delta G^\ddagger = \Delta G + \frac{\lambda}{\ln 2} \ln \{1 + \exp[-(\Delta G \ln 2)/\lambda]\} \quad (1b)$$

where  $k_d$  and  $k_{-d}$  are the rate constants for forward and reverse diffusion in solution,  $k_{\text{en}}^0$  is a pre-exponential factor,  $\lambda$  is the reorganizational energy and  $\Delta G = (E_{\text{Acceptor}}^{00} - E_T)$ . By using parameters derived from a previous study<sup>13</sup> ( $k_d = 1.0 \times 10^{10}$

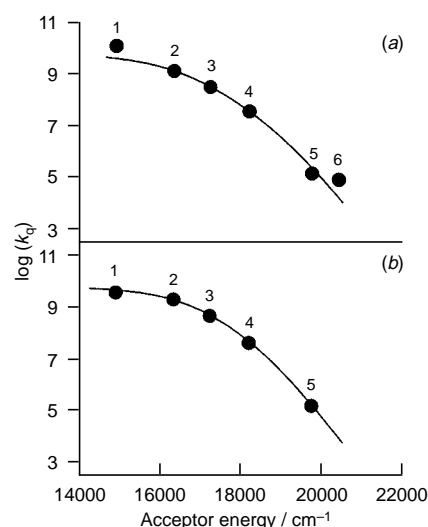


Fig. 2 Sandros plots for (a) **P0** and (b) **1**. Points are experimentally determined rate constants with various triplet energy acceptors,<sup>¶</sup> and fitted lines are calculated as explained in the text.

**Table 1** Photophysical properties of model compound **1** and polymer **P0**<sup>a</sup>

Compound	$\tau_{\text{fl}}/\text{ns}$	$\Phi_{\text{fl}}^b$	$E_{\text{S}}/\text{cm}^{-1}$ (kcal mol <sup>-1</sup> )	$\lambda^c/\text{cm}^{-1}$	$E_{\text{T}}^c/\text{cm}^{-1}$ (kcal mol <sup>-1</sup> )	$\Phi_{\text{T}}$	$10^{-8} k_{\text{p}}/\text{s}^{-1}$	$10^{-8} k_{\text{nr}}/\text{s}^{-1}$	$10^{-8} k_{\text{isc}}/\text{s}^{-1}$
<b>1</b>	1.15	0.90	24 200 (69.2)	1000	17 700 (50.6)	$\leq 0.10^{d,e}$	7.8	$\leq 0.87$	$\leq 0.87$
<b>P0</b>	0.53	0.60 <sup>f</sup>	23 100 (66.0)	1200	17 500 (51.2)	0.12 <sup>e</sup>	11	5.7	2.3

<sup>a</sup> THF solutions at 298 K. <sup>b</sup> Fluorescence quantum yields determined using 9,10-dicyanoanthracene (DCA) and perylene as actinometers ( $\Phi_{\text{DCA}}^{\text{EtOH}} = 0.89$ ,  $\Phi_{\text{Perylene}}^{\text{EtOH}} = 0.89$ ). <sup>c</sup> Values based on a Marcus fit to quenching data as described in text. <sup>d</sup> Upper limit. <sup>e</sup> See footnote §. <sup>f</sup> See footnote \*\*.

$\text{M}^{-1} \text{s}^{-1}$ ,  $k'_{-d} = 2.0 \times 10^{10} \text{M}^{-1} \text{s}^{-1}$ ,  $k_{\text{en}}^0 = 5 \times 10^{10} \text{s}^{-1}$ ), best fits of the Sandros plots for **1** and **P0** were attained by varying  $\lambda$  and  $E_{\text{T}}$  in eqn. (1). These values are listed in Table 1, and plots of the fitted quenching data are shown in Fig. 2. Surprisingly,  $E_{\text{T}}$  for the model compound and the polymer is the same within experimental error ( $\pm 5\%$ ).

A second objective of the present study was to determine the triplet yields ( $\Phi_{\text{T}}$ ) for **1** and **P0**. However, in order to determine this parameter it is necessary to know the fluorescence quantum yields ( $\Phi_{\text{fl}}$ ). Therefore, fluorescence spectra of **1** and **P0** were obtained [ $\lambda_{\text{max}}(\mathbf{1}) = 413 \text{ nm}$ ,  $\lambda_{\text{max}}(\mathbf{P0}) = 433 \text{ nm}$ ], and  $\Phi_{\text{fl}}$  values were determined from the integrated fluorescence spectra (Table 1).<sup>14</sup> Fluorescence lifetimes ( $\tau_{\text{fl}}$ ) for **1** and **P0** are also listed in Table 1. Next, the triplet yield ( $\Phi_{\text{T}}$ ) for **P0** was determined by time-resolved thermal lensing (TRTL).<sup>15,16</sup> In this method,  $\Phi_{\text{T}}$  is derived from the ratio of the 'slow' and 'total' heat deposition signals ( $U_{\text{slow}}$  and  $U_{\text{total}}$ , respectively) after correcting for the energy of the pump laser ( $E_{\text{hv}}$ ) and the fluorescence quantum yield,

$$U_{\text{slow}}/U_{\text{total}} = E_{\text{T}}\Phi_{\text{T}}/(E_{\text{hv}} - E_{\text{S}}\Phi_{\text{fl}}) \quad (2)$$

where  $E_{\text{S}}$  is the energy of the singlet state (Table 1). The  $U_{\text{slow}}/U_{\text{total}}$  ratio of a vacuum degassed THF solution of **P0** was determined to be 0.15, and the triplet yield calculated with eqn. (2) is listed in Table 1.\*\* Unfortunately,  $\Phi_{\text{T}}$  for **1** could not be measured with TRTL owing to strong triplet–triplet absorption at the thermal lensing probe wavelength (633 nm). However, the experimentally derived  $\Phi_{\text{fl}}$  for **1** (0.90) provides an upper limit for  $\Phi_{\text{T}}$  ( $\leq 0.1$ ).\*\*

The photophysical characterization of the singlet and triplet states of **1** and **P0** allows us to assess the effect that the extended  $\pi$ -conjugation in the polymer has on its triplet state properties. It is surprising that  $E_{\text{T}}$  is the same for the polymer and model compound. However, the more important parameter is the singlet–triplet splitting energy ( $E_{\text{ST}}$ ), which is defined as  $E_{\text{S}} - E_{\text{T}}$ . Interestingly,  $E_{\text{ST}}$  is lower by almost  $1000 \text{ cm}^{-1}$  in **P0** compared to **1** ( $5600 \text{ vs. } 6500 \text{ cm}^{-1}$ ). The lower  $E_{\text{ST}}$  indicates that the electron exchange energy is smaller in the polymer,<sup>17</sup> which presumably is due to the fact that the delocalization of the LUMO (and/or the HOMO) for the  $\pi \rightarrow \pi^*$  transition is greater in **P0** than in **1**. Moreover, it is of interest that the experimentally determined  $E_{\text{ST}}$  for **P0** is comparable to values calculated using semi-empirical MO theory.<sup>18,19</sup>

The rates of radiative decay, non-radiative decay and intersystem crossing ( $k_{\text{r}}$ ,  $k_{\text{nr}}$  and  $k_{\text{isc}}$ , respectively) for the  $^1\pi, \pi^*$  state of **1** and **P0** have been determined by using the experimental  $\Phi_{\text{fl}}$ ,  $\tau_{\text{fl}}$  and  $\Phi_{\text{T}}$  values (Table 1). First, note that  $k_{\text{r}}$  is comparable in **1** and **P0**. This is not surprising given that the orbital basis and transition dipole for radiative transitions are similar in the model and the polymer. By contrast,  $k_{\text{nr}}$  and  $k_{\text{isc}}$  are significantly larger in **P0** than in **1**. The larger non-radiative decay rate in the polymer may arise from 'defect sites' (e.g. unreacted end groups and/or defects in the polymer backbone) that quench the  $^1\pi, \pi^*$  exciton. The larger (effective) intersystem crossing rate in **P0** may arise from the smaller  $E_{\text{ST}}$  value.

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## Notes and References

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‡ **P0** was synthesized via a Pd-mediated cross-coupling reaction of the polymer subunits (see ref. 7 for complete synthetic details). *Selected characterization*: **P0**: GPC ( $\text{CHCl}_3$ , polystyrene standards)  $M_n = 65.2 \text{ kD}$ ,  $M_w = 209.2 \text{ kD}$  (PDI = 3.2). <sup>1</sup>H NMR ( $\text{CDCl}_3$ )  $\delta$  0.87 (bt, t), 1.24 (br, s), 1.56 (br, m), 1.84 (br, m), 4.04 (br, t), 6.92 (br, s), 7.0 (br, s), 7.62 (br, s). <sup>13</sup>C NMR ( $\text{CDCl}_3$ )  $\delta$  14.1, 22.7, 26.1, 29.4, 29.7, 32.0, 69.7, 87.0, 94.8, 114.0, 116.9, 122.8, 126.8, 132.1, 140.0, 153.7. **1**: <sup>1</sup>H NMR ( $\text{CDCl}_3$ )  $\delta$  3.79 (s, 6 H, OCH<sub>3</sub>), 3.89 (s, 6 H, OCH<sub>3</sub>), 6.86 (m, 4 H, phenyl), 7.08 (d, 2 H, phenyl), 7.63 (d, 8 H, biphenyl). <sup>13</sup>C NMR ( $\text{CDCl}_3$ )  $\delta$  55.8, 56.5, 86.7, 93.2, 112.2, 112.9, 115.8, 118.0, 122.7, 126.7, 132.1, 139.9, 153.2, 154.5.

§ In ref. 6 we reported the triplet lifetime for a sample of **P0** in Ar-purged THF ( $\tau_{\text{T}} = 4 \mu\text{s}$ ). In that study the lifetime was suppressed due to quenching by residual oxygen. The triplet lifetimes reported herein are for samples that were freeze–pump–thaw degassed five times and sealed at  $10^{-5}$  Torr. It is conceivable that even under these conditions the triplet lifetime is suppressed by oxygen quenching.

¶ Triplet acceptors (see Fig. 2): 1, anthracene; 2, [4,4'-bis(carboethoxy)-2,2'-bipyridyl]Re(CO)<sub>3</sub>Cl; 3, *trans*-stilbene; 4, (2,2'-bipyridyl)Re(CO)<sub>3</sub>Cl; 5, biacetyl; 6, *p*-terphenyl.

|| In ref. 6 we reported the fluorescence quantum yield for **P0** in Ar-purged THF ( $\Phi_{\text{fl}} = 0.28$ ). The sample of **P0** used in that study was structurally equivalent to that used in the present study, but it had a substantially lower molecular weight ( $M_n = 13.5 \text{ kD}$ ,  $M_w = 37.4 \text{ kD}$ ). The higher molecular weight of the sample of **P0** used in the present study accounts for the larger fluorescence yield reported herein.

\*\* The  $\Phi_{\text{T}}$  values for **P0** and **1** were confirmed using photoacoustic calorimetry [ $\Phi_{\text{T}}(\mathbf{P0}) = 0.12$ ,  $\Phi_{\text{T}}(\mathbf{1}) = 0.05$ ].

- L. J. Rothberg, M. Yan, F. Papadimitrakopoulos, M. E. Galvin, E. W. Kwock and T. M. Mille, *Synth. Met.*, 1996, **80**, 41 and references therein.
- B. Xu and S. Holdcroft, *Adv. Mater.*, 1994, **6**, 325.
- M. S. A. Abdou, F. P. Orfino, Y. Son and S. Holdcroft, *J. Am. Chem. Soc.*, 1997, **119**, 4518.
- R. D. Scurlock, B. Wang, P. R. Ogilby, J. R. Sheats and R. L. Clough, *J. Am. Chem. Soc.*, 1995, **117**, 10 194.
- H. F. Wittman, R. H. Friend, M. S. Khan and J. Lewis, *J. Chem. Phys.*, 1994, **101**, 2693.
- K. D. Ley, C. E. Whittle, M. D. Bartberger and K. S. Schanze, *J. Am. Chem. Soc.*, 1997, **119**, 3423.
- K. D. Ley and K. S. Schanze, *Coord. Chem. Rev.*, in press.
- S. L. Murov, I. Carmichael and G. L. Hug, *Handbook of Photochemistry*, Marcel Dekker, New York, 1993, p. 54.
- L. A. Worl, R. Duesing, P. Chen, L. D. Ciana and T. J. Meyer, *J. Chem. Soc., Dalton Trans.*, 1991, 849.
- K. Sandros, *Acta Chem. Scand.*, 1964, **18**, 2355.
- V. Balzani and F. Bolletta, *J. Am. Chem. Soc.*, 1978, **100**, 7404.
- V. Balzani, F. Bolletta and F. Scandola, *J. Am. Chem. Soc.*, 1980, **102**, 2152.
- Y. Wang and K. S. Schanze, *Inorg. Chem.*, 1994, **33**, 1354.
- G. A. Crosby and J. N. Demas, *J. Phys. Chem.*, 1971, **75**, 991.
- S. E. Braslavsky and G. E. Heibel, *Chem. Rev.*, 1992, **92**, 1381.
- R. T. Cambron and J. M. Harris, *J. Phys. Chem.*, 1993, **97**, 13 598.
- S. P. McGlynn, F. J. Smith and G. Cilento, *Photochem. Photobiol.*, 1964, **3**, 269.
- J. W. Blatchford, T. L. Gustafson and A. J. Epstein, *J. Chem. Phys.*, 1996, **105**, 9214.
- D. Beljonne, H. F. Wittmann, A. Köhler, S. Graham, M. Yonnis, J. Lewis, P. R. Raithby, M. S. Khan, R. H. Friend and J. L. Brédas, *J. Chem. Phys.*, 1996, **105**, 3868.

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